

## When Is a Measurement Reversible?

Information Gain and Disturbance in Quantum Measurements (revisited)  
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these slides are available for download at [tinyurl.com/BDW-aqis16](http://tinyurl.com/BDW-aqis16)



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- ✓ mathematically, a measurement process is represented by a **completely positive instrument** (Ozawa, 1984), namely, a family  $\mathfrak{M} = \{\mathcal{E}_m : m \in \mathcal{M}\}$  of completely positive maps  $\mathcal{E}_m$ , indexed by the outcome set  $\mathcal{M} = \{m\}$ , and such that  $\sum_m \mathcal{E}_m$  is trace-preserving (completeness relation)



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- ✓ upon input  $\rho$ , the outcome  $m$  is obtained with probability  $p(m) = \text{Tr}[\mathcal{E}_m(\rho)]$  and the corresponding post-measurement state is  $\sigma_m = \mathcal{E}_m(\rho)/p(m)$

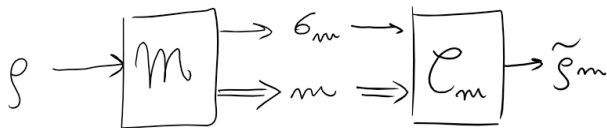
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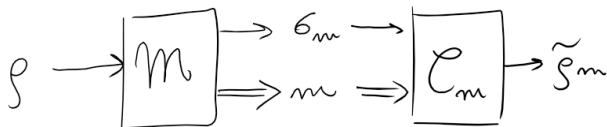
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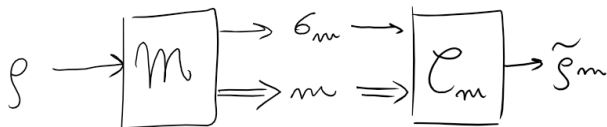
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- ✓ **Fact**. Among all measurements with the same outcome statistics, **efficient measurements** (i.e.,  $\mathcal{E}_m(\cdot) = E_m \cdot E_m^\dagger$ ) are the least disturbing one: we will hence focus on these only



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- ✓ tempting analogy with **work and heat in thermodynamics**

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## Stronger than any other (currently known) recoverability results:

- 1 the lemma holds also for positive (not necessarily completely positive)  $\Phi$
- 2 the explicit form of correction  $\Psi$  is known
- 3  $D(\rho \| \sigma) \geq D_{\text{M}}(\rho \| \sigma) \geq -\log F(\rho, \sigma)$ , where  $F(\rho, \sigma) \triangleq \|\sqrt{\rho}\sqrt{\sigma}\|_1^2$

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- ✓ in other words, the channel  $\mathcal{M}$  is subunital and the Technical Lemma can be applied



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- ✓ hence, the “extrinsic noise” measures the irreversibility of efficient measurements (as does entropy for adiabatic processes)
- ✓ indeed, in [K. Jacobs, PRA **80**, 012322 (2009)] the inequality  $H(M) - I_G(\rho, \mathfrak{M}) \geq 0$  is interpreted as the second law for quantum measurements; the above Theorem considerably strengthens this



Thank You

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