

# Manipulating quantum information of two trapped ions by a single-step operation

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Based on the exact conditional quantum dynamics for a two-ion system, here we propose an efficient *single-step* scheme for coherently manipulating quantum information of two trapped cold ions beyond the Lamb-Dicke limit by using a pair of synchronous laser pulses.

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## I. INTRODUCTION

The entanglement between different qubits has recently become a focus of activity in quantum physics[1], because of experiments on non-local features of quantum mechanics and the development of quantum information physics. In the past few years, several key features of the proposal in [2], including the production of entangled states and the implementation of quantum controlled operations between a pair of trapped ions, have already been experimentally demonstrated[4–6]. Meanwhile, several alternative theoretical schemes[7–10] have also been developed for overcoming various difficulties in realizing a practical ion-trap quantum information processor.

## II. MODEL

In this work, we propose an effective scheme for realizing the communication and logic operations between different trapped ions by a *single-step* operation, performed by using different laser beams synchronously. Following Sørensen and Mølmer [10], the Hamiltonian describing two trapped cold ions driven by a pair of synchronous laser beams is

$$\hat{H}(t) = \hbar\nu \left( \hat{a}^\dagger \hat{a} + \frac{1}{2} \right) + \hbar\omega_0 \sum_{j=1}^2 \frac{\hat{\sigma}_{z,j}}{2} + \frac{\hbar}{2} \sum_{j=1}^2 \Omega_j \{ \hat{\sigma}_{+,j} \exp\{i[\eta_j(\hat{a}^\dagger + \hat{a}) - \omega_j t - \phi_j]\} + H.c. \}. \quad (1)$$

Under the rotating wave approximation, this system is exactly solvable *even if the usual Lamb-Dicke approximation is not made*. For brevity, the expressions for the solutions are omitted here and will be presented elsewhere.

## III. RESULTS

Based on the exact solutions of the rotating wave approximation to the Hamiltonian (1), here we focus on how to realize in *one step* both two-qubit controlled operations and two-qubit entangled states between ion 1 and 2 beyond the Lamb-Dicke limit. This is achieved by properly setting up the controllable experimental parameters, e.g., the Lamb-Dicke parameters  $\eta_j$ , the carrier Rabi frequencies  $\Omega_j$ , the frequencies  $\omega_j$  ( $j = 1, 2$ ) and duration of the applied synchronous pulses.

1. The  $\hat{C}^Z$  gate

$$\hat{C}_{12}^Z = |g_1\rangle|g_2\rangle\langle g_1|\langle g_2| + |g_1\rangle|e_2\rangle\langle g_1|\langle e_2| + |e_1\rangle|g_2\rangle\langle e_1|\langle g_2| - |e_1\rangle|e_2\rangle\langle e_1|\langle e_2|, \quad (2)$$

can be realized directly by using a pair of red-sideband pulses.

2. if a resonant pulse is applied to ion 2 and a off-resonant pulse is simultaneously applied to ion 1, we find that the two-qubit controlled operation

$$\hat{C}_{12} = |g_1\rangle|g_2\rangle\langle g_1|\langle g_2| + |g_1\rangle|e_2\rangle\langle g_1|\langle e_2| \pm i e^{-i\phi_2} |e_1\rangle|g_2\rangle\langle e_1|\langle e_2| \pm i e^{i\phi_2} |e_1\rangle|e_2\rangle\langle e_1|\langle g_2|, \quad (3)$$

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can be implemented directly. These gates are equivalent to the exact CNOT gate, except for a local rotation.

### 3. Two-qubit entangled states

$$|\psi_{12}^-\rangle = U |g_1\rangle|e_2\rangle - V |e_1\rangle|g_2\rangle, \quad |\psi_{12}^+\rangle = V |g_1\rangle|e_2\rangle + U |e_1\rangle|g_2\rangle, \quad (4)$$

and

$$|\phi_{12}^-\rangle = U |g_1\rangle|g_2\rangle - V |e_1\rangle|e_2\rangle, \quad |\phi_{12}^+\rangle = V |g_1\rangle|g_2\rangle + U |e_1\rangle|e_2\rangle, \quad (5)$$

can be generated from the dynamical evolutions of the non-entangled initial states  $|g_1\rangle|e_2\rangle$ ,  $|e_1\rangle|g_2\rangle$  and  $|g_1\rangle|g_2\rangle$ ,  $|e_1\rangle|e_2\rangle$ , respectively. These entangled states become the relevant two-qubit maximally entangled states, i.e., EPR states:

$$|\Phi_{12}^\pm\rangle = \frac{1}{\sqrt{2}} (|g_1\rangle|g_2\rangle \pm |e_1\rangle|e_2\rangle), \quad |\Psi_{12}^\pm\rangle = \frac{1}{\sqrt{2}} (|g_1\rangle|e_2\rangle \pm |e_1\rangle|g_2\rangle), \quad (6)$$

if the experimental parameters are set up properly.

4. The EPR states, e.g.,  $|\Phi_{12}^-\rangle$  and  $|\Psi_{12}^+\rangle$ , can also be generated by sequentially using the a single-qubit rotation and the controlled operation introduced above,

$$\left\{ \begin{array}{l} |m\rangle|g_1\rangle|g_2\rangle \xrightarrow{\hat{r}_1} |m\rangle \otimes \frac{1}{\sqrt{2}} (|g_1\rangle|g_2\rangle - i e^{-i\varphi_1} |e_1\rangle|g_2\rangle) \xrightarrow{\hat{C}_{12}} |m\rangle \otimes |\Phi_{12}^-\rangle, \\ |m\rangle|g_1\rangle|g_2\rangle \xrightarrow{\hat{r}_1} |m\rangle \otimes \frac{1}{\sqrt{2}} (|g_1\rangle|g_2\rangle, - i e^{-i\varphi_1} |e_1\rangle|g_2\rangle) \xrightarrow{\hat{C}'_{12}} |m\rangle \otimes |\Psi_{12}^+\rangle. \end{array} \right. \quad (7)$$

## IV. CONCLUSIONS

Based on the exact conditional quantum dynamics for two trapped ions driven by a pair of synchronous laser beams, we have shown that, under certain conditions, the quantum controlled gates and entanglement between a pair of trapped ions can be realized deterministically by *only a single-step operation*. The CM mode of ions always remains in its initial quantum state after the operation.

Compared to other approaches for coherently manipulating a pair of ions, the present scheme has some advantages. First, compared to the the previous multi-pulse schemes for two-qubit controlled gates (see, e.g., [2, 8, 9]), the present single-step operational scheme may be easily tested experimentally, as the undesired dynamical phase evolution [11] related to the delay time between different operations is avoided completely. It may be more important that the present scheme can be used to describe a laser-ion coupling of arbitrary strength, while previous schemes (see, e.g., [2, 10]) based on the LD approximation are only effective in the weak-coupling regime.

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